

Hilbert Space / Observables

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Note Title

- A set of all vectors $|\alpha\rangle$ that are closed for the standard vector operations are called "vector space"
- A similar set of all wave functions defined over an interval $a \leq x \leq b$ is called a "Hilbert Space" for that particular interval.

- wave functions $\rightarrow$ square-integrable
 $\Leftrightarrow \int_a^b |f(x)|^2 dx < \infty$

- Inner product of two functions $f(x)$ & $g(x)$

$$\langle f|g \rangle \equiv \int_a^b f(x)^* g(x) dx$$

Compare for vectors

$$\langle \alpha|\beta \rangle = a_1^* b_1 + a_2^* b_2 + \dots$$

- $\langle g|f \rangle = \langle f|g \rangle^*$
- $\langle f|f \rangle = \int_a^b |f(x)|^2 dx \geq 0$

- A set of functions is orthonormal
if $\langle f_m|f_n \rangle = \delta_{mn}$,

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- A set of functions is complete if any function in the Hilbert space can be expressed as a linear combination of them:

$$f(x) = \sum_{n=1}^{\infty} c_n f_n(x)$$

$$\Rightarrow c_n = \langle f_n | f \rangle$$

$$\Gamma \equiv \int_a^b f_n(x)^* f(x) dx$$

- Ex.** ① The set of energy eigenfunctions of the infinite square potential well over $[0, a]$.
 ② The set of energy eigenfunctions of the harmonic potential well over $[-\infty, \infty]$

Observables

- An "observable" operator Q always gives a real value for the measurement.
 That is $\langle Q \rangle = \langle Q \rangle^*$
- This leads to the requirement that any observable operator should be hermitian (See the text book for proof)
- What does "hermitian" operator mean?

From Appendix A.6

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- A matrix T is called hermitian if

$$T^\dagger \equiv (T^*)^T = T$$

Ex)

$$T = \begin{pmatrix} 1 & i \\ i & 2 \end{pmatrix} \Rightarrow T^\dagger = \begin{pmatrix} 1 & -i \\ -i & 2 \end{pmatrix}$$

So this is not hermitian

$$T = \begin{pmatrix} 1 & i \\ -i & 2 \end{pmatrix} \Rightarrow T^\dagger = \begin{pmatrix} 1 & i \\ -i & 2 \end{pmatrix}$$

$\Rightarrow$ this is hermitian

$$T = \begin{pmatrix} 1 & i \\ 1 & 2 \end{pmatrix} \Rightarrow T^\dagger = \begin{pmatrix} 1 & 1 \\ -i & 2 \end{pmatrix}$$

$\Rightarrow$ this is not hermitian

↑ Here $T^\dagger$ is called the hermitian conjugate of T

- For two matrices S & T

$$(S \cdot T)^\dagger = T^\dagger S^\dagger, \text{ where}$$

$T^\dagger$ is called hermitian conjugate of T , and $(T^\dagger)^\dagger = T$

- For two vectors and a matrix

$$\vec{a}^\dagger T \cdot \vec{b} = (T^\dagger \vec{a})^\dagger \vec{b}$$

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• In Dirac notations,

$$\vec{a}^\dagger T \vec{b} \Leftrightarrow \langle \alpha | T | \beta \rangle$$

$$(T^\dagger \vec{a})^\dagger \vec{b} \Leftrightarrow \langle T^\dagger \alpha | \beta \rangle$$

So

$$\langle T^\dagger \alpha | \beta \rangle = \langle \alpha | T | \beta \rangle$$

for any matrix (or operator) T

If T is a hermitian matrix
 $T^\dagger = T$.

In other words $\langle T \alpha | \beta \rangle = \langle \alpha | T | \beta \rangle$

Ex. $|\alpha\rangle = \begin{pmatrix} \bar{z} \\ 1 \end{pmatrix}, \quad |\beta\rangle = \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix}$

check for $T = \begin{pmatrix} 0 & \bar{z} \\ 1 & 0 \end{pmatrix},$

(non hermitian)

$$\langle \alpha | T | \beta \rangle$$

$$= (-\bar{z} \ 1) \begin{pmatrix} 0 & \bar{z} \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} = (-\bar{z} \ 1) \begin{pmatrix} -1 \\ 1 \end{pmatrix}$$

$$= \bar{z} + 1$$

$$\langle T \alpha | \beta \rangle = \left[\begin{pmatrix} 0 & \bar{z} \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \bar{z} \\ 1 \end{pmatrix} \right]^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix}$$

$$= \begin{pmatrix} \bar{z} \\ \bar{z} \end{pmatrix}^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} = (-\bar{z}, -\bar{z}) \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} = -\bar{z} + 1$$

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So $\langle \alpha | T \beta \rangle \neq \langle T \alpha | \beta \rangle$
for this case.

$$\begin{aligned}
 \bullet \text{ But if we do } & \langle T^\dagger \alpha | \beta \rangle \\
 &= \left[\begin{pmatrix} 0 & 1 \\ -\bar{z} & 0 \end{pmatrix} \begin{pmatrix} \bar{z} \\ 1 \end{pmatrix} \right]^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} \\
 &= \begin{bmatrix} 1 \\ 1 \end{bmatrix}^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} = (1 \ 1) \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} \\
 &= 1 + \bar{z} = \langle \alpha | T \beta \rangle.
 \end{aligned}$$

• Now if we check for $T = \begin{pmatrix} 0 & \bar{z} \\ -\bar{z} & 1 \end{pmatrix}$
which is hermitian ($\because T^\dagger = T$)

$$\begin{aligned}
 \langle \alpha | T \beta \rangle &= (-\bar{z} \ 1) \begin{pmatrix} 0 & \bar{z} \\ -\bar{z} & 1 \end{pmatrix} \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} \\
 &= (-\bar{z} \ 1) \begin{pmatrix} -1 \\ 0 \end{pmatrix} = \underline{\bar{z}}
 \end{aligned}$$

$$\begin{aligned}
 \langle T \alpha | \beta \rangle &= \left[\begin{pmatrix} 0 & \bar{z} \\ -\bar{z} & 1 \end{pmatrix} \begin{pmatrix} \bar{z} \\ 1 \end{pmatrix} \right]^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} \\
 &= \begin{bmatrix} \bar{z} \\ 2 \end{bmatrix}^\dagger \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} = (-\bar{z} \ 2) \begin{pmatrix} 1 \\ \bar{z} \end{pmatrix} \\
 &= -\bar{z} + 2\bar{z} = \underline{\bar{z}}
 \end{aligned}$$

Thus $\langle \alpha | T \beta \rangle = \langle T \alpha | \beta \rangle$
as it should be for a hermitian matrix

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- * The most famous property of hermitian operator is that its eigenvalues are always real.
This is also why all observables are represented by hermitian operators

$$\begin{array}{ccc} T | \alpha \rangle = \lambda | \alpha \rangle & & \\ \uparrow & \swarrow & \\ \text{hermitian operator} & & \text{eigenvalue (real)} \end{array}$$

- * Example of a non-matrix hermitian operator: $\hat{p} = \frac{\hbar}{i} \frac{\partial}{\partial x}$

Check if $\langle f | \hat{p} g \rangle = \langle \hat{p} f | g \rangle$

$$\begin{aligned} \langle f | \hat{p} g \rangle &= \frac{\hbar}{i} \int_{-\infty}^{\infty} f^*(x) \frac{d}{dx} g(x) dx \\ &= \frac{\hbar}{i} \left[\cancel{f^*(x) g(x)} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{d}{dx} f^*(x) g(x) dx \right] \\ &= \int_{-\infty}^{\infty} \left(\frac{\hbar}{i} \frac{d}{dx} f(x) \right)^* g(x) dx \\ &= \int_{-\infty}^{\infty} (\hat{p} f(x))^* g(x) dx = \langle \hat{p} f | g \rangle \end{aligned}$$

∴ $\hat{p}$ is an hermitian operator

Determinate States

⑦

- * Measurement of an observable Q on a state $|\Psi\rangle$ is certain to yield the same value of q ,
 $|\Psi\rangle$ is called a "determinate state", and this process can be mathematically expressed as

$$\hat{Q}|\Psi\rangle = q|\Psi\rangle$$

This type of equation is called eigenvalue equation. $|\Psi\rangle$ is an "eigenfunction" and " q " is the corresponding "eigenvalue".

- * Determinate states of an observable $\hat{Q}$ are mathematically eigenfunctions of $\hat{Q}$
- * Collection of all eigenvalues are called "spectrum"
- * If two or more linearly independent eigenfunctions share the same eigenvalue the spectrum is said to be "degenerate".

- * Ex. $\hat{H}|\Psi\rangle = E|\Psi\rangle$
 E : eigenvalues of the Hamiltonian $\hat{H}$
 $|\Psi\rangle$: stationary states = eigen ftns of $\hat{H}$.

Q

Ex Eigen value problem with matrices

$$T = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

Find its eigenvalues and eigenvectors.

$$T|\alpha\rangle = \lambda|\alpha\rangle, \quad |\alpha\rangle = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$$

$$\Rightarrow (T - \lambda I)|\alpha\rangle = 0$$

$$\Rightarrow \text{For } |\alpha\rangle \text{ to be non-zero, } \det(T - \lambda I) = 0$$

$$\Rightarrow \det \left[\begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} - \lambda \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right] = \det \begin{pmatrix} -\lambda & i \\ i & -\lambda \end{pmatrix}$$

$$= \lambda^2 - i^2 = (\lambda - i)(\lambda + i) = 0$$

$$\therefore \text{eigenvalues } \underline{\lambda = i, -i}$$

eigen functions ?

$$\text{For } \lambda = i, \quad \begin{pmatrix} -i & i \\ i & -i \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = 0$$

$$\Rightarrow a_1 - a_2 = 0 \Rightarrow a_1 = a_2$$

$$\therefore |\alpha\rangle = \begin{pmatrix} 1 \\ 1 \end{pmatrix} a_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

(normalized)

$$\text{For } \lambda = -i, \quad a_1 + a_2 = 0 \Rightarrow a_2 = -a_1$$

$$|\alpha\rangle = \begin{pmatrix} 1 \\ -1 \end{pmatrix} a_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

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Ex

Now non-matrix operator

$$\hat{Q} \equiv \bar{L} \frac{d}{d\phi}, \quad \phi \text{ is the polar coordinate in 2D.}$$

Find the eigenvalues and eigenfunctions, check if it is hermitian

$$\hat{Q} |\psi\rangle = \lambda |\psi\rangle$$

$$\text{with } |\psi\rangle = \psi(\phi)$$

$$\bar{L} \frac{d}{d\phi} \psi(\phi) = \lambda \psi(\phi) \Rightarrow \frac{\bar{L}}{\lambda} \frac{d\psi(\phi)}{d\psi} = d\phi$$

$$\Rightarrow \ln \psi(\phi) = \frac{-i\lambda}{\bar{L}} \phi + \text{const}$$

$$\Rightarrow \psi(\phi) = A e^{-i\lambda \phi / \bar{L}}$$

$$\psi(0) = \psi(2\pi) \Rightarrow e^{-i\lambda 2\pi / \bar{L}} = e^0 = 1$$

$$\Rightarrow -i\lambda 2\pi / \bar{L} = 2n\pi i \Rightarrow \lambda = -n$$

$$, n = 0, \pm 1, \pm 2, \dots$$

So eigenvalues, $\lambda =$ all integers
eigenfunctions, $\psi(\phi) = A e^{-i\lambda \phi / \bar{L}}$

$$\langle f | \hat{Q} g \rangle =$$

$$\int_0^{2\pi} f^*(\phi) \left(\bar{L} \frac{d}{d\phi} \right) g(\phi) d\phi = \bar{L} \left[f^*(\phi) g(\phi) \Big|_0^{2\pi} - \int_0^{2\pi} \frac{d f^*(\phi)}{d\phi} g(\phi) d\phi \right]$$

$$= \int_0^{2\pi} \left(\bar{L} \frac{d f(\phi)}{d\phi} \right)^* g(\phi) d\phi$$

$$= \langle \hat{Q} f | g \rangle \quad \therefore \text{hermitian}$$